

CUF-2013-1sem

22/03/2016

Q. 40

(R) Q

Q1. a)  $r > R$

$$E = \frac{Q_{enc}}{4\pi r^2 \epsilon_0} = \frac{Q}{4\pi r^2 \epsilon_0}$$

$$Q_{enc} = Q$$

$r \leq R$

$$E = \frac{Q_{enc}}{4\pi r^2 \epsilon_0} = \frac{rQ}{4\pi R^3 \epsilon_0}$$

$$Q_{enc} = \frac{4\pi r^3}{3} \cdot \frac{Q}{4\pi R^3} = \frac{r^3}{R^3} Q$$

b)  $dF = \frac{q \cdot dV \cdot Q}{4\pi r^2 \epsilon_0 V}$  ?

c)  $F = \frac{qQ}{2} \int_0^R \frac{4\pi r^2 dr}{4\pi r^2 \epsilon_0 \cdot 4\pi R^3} = \frac{qQ}{2\epsilon_0 R^2}$  ?

Q2

a)  $E = \frac{\nabla \cdot \hat{z}}{\epsilon_0} = \frac{q_0 \sin(\omega t)}{2\pi a^2 \epsilon_0} \hat{z}$  ✓

b)  $\nabla \times E = -\frac{\partial E_z}{\partial t} \hat{\phi} = 0 \Rightarrow \int H \cdot dl = \frac{\partial}{\partial t} \int D \cdot ds \Rightarrow B = \mu_0 \epsilon_0 \frac{\partial E}{\partial t} \cdot \frac{\pi a^2}{2\pi}$

$$B = \frac{\mu_0 q_0 \omega \sin(\omega t) \cdot \pi a^2}{2\pi a^2}$$

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$$c) S = \frac{E \cdot B}{\mu_0} \Rightarrow \frac{E_z \cdot B_\phi}{\mu_0} (-\hat{r}) = \left( \frac{q_0 \sin(\omega t)}{\hat{n} a^2 \epsilon_0} \right) \cdot \left( \frac{q_0 \cos(\omega t) \cdot \omega r}{2 \hat{n} a^2} \right) (-\hat{r}) \quad \checkmark$$

$$E \times B = \begin{vmatrix} \hat{r} & \hat{\theta} & \hat{\phi} \\ E_r & E_\theta & E_\phi \\ B_r & B_\theta & B_\phi \end{vmatrix} = -E_z B_\phi \hat{r}$$

$$S = \frac{q_0^2 \omega r \sin(\omega t) \cos(\omega t)}{2 \hat{n}^2 a^4 \epsilon_0} (-\hat{r}) = \frac{q_0^2 \omega r \sin(2\omega t)}{(2 \hat{n} a)^2 \epsilon_0} \quad \checkmark$$

$$d) \frac{\omega a}{c} \ll 1 \quad \mu = \frac{1}{2} \left( \frac{\epsilon_0 q_0^2 \sin^2(\omega t)}{(\hat{n} c^2 \epsilon_0)^2} + \mu_0 \frac{\omega^2 r^2 \cos^2(\omega t)}{(2 \hat{n} a)^2 \epsilon_0} \right) = \frac{q_0^2 \mu_0}{2 \hat{n}^2 a^4} \left[ \frac{c^2 \sin^2(\omega t)}{4} + \frac{\omega^2 r^2 \cos^2(\omega t)}{4} \right]$$

$$\frac{\partial \mu}{\partial t} = \frac{q_0^2 \mu_0}{2 \hat{n}^2 a^4} \left[ \underbrace{c^2 \cdot 2 \sin(\omega t) \cdot \cos(\omega t) \cdot \omega}_{\frac{1}{2} \sin(2\omega t)} + \underbrace{\omega^2 r^2 \cdot 2 \cos(\omega t) \cdot (-\sin(\omega t) \cdot \omega)}_{\frac{1}{2} \sin(2\omega t)} \right]$$

$$\nabla S = \frac{1}{r} \frac{\partial(r S_\phi)}{\partial r} = \frac{1}{r} \cdot \frac{q_0^2 \omega r \sin(2\omega t) \cos(\omega t)}{\hat{n}^2 a^4 \epsilon_0} = \frac{q_0^2 \omega \sin(2\omega t)}{2 \hat{n}^2 a^4 \epsilon_0}$$

$$\nabla S + \frac{\partial \mu}{\partial t} = 0 \Rightarrow \frac{q_0^2 \omega \sin(2\omega t)}{2 \hat{n}^2 a^4 \epsilon_0} + \frac{q_0^2 \mu_0 \sin(2\omega t) \omega}{2 \hat{n}^2 a^4} \left[ \frac{c^2}{4} - \frac{\omega^2 r^2}{4} \right]$$

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$$\frac{q_0^2 \omega \sin(2\omega t)}{2\hbar c^4 \epsilon_0} + \frac{q_0 \omega \sin(2\omega t)}{2\hbar c^4 \epsilon_0} \left[ 1 - \left( \frac{\omega r}{c} \right)^2 \right] = 0 \quad \checkmark$$

Q3.

$$a) |\langle \psi | \psi \rangle|^2 = 1 \Rightarrow \int_{-\infty}^{\infty} |\psi(x)|^2 dx = 1 \Rightarrow \int_{-\frac{L}{2}}^{\frac{L}{2}} A^2 \cos^2\left(\frac{3\pi x}{L}\right) dx = 1$$

$$A^2 \left[ \int_{-\frac{L}{2}}^{\frac{L}{2}} \frac{1}{2} + \frac{\cos\left(\frac{6\pi x}{L}\right)}{2} dx \right] = 1 \Rightarrow A^2 \cdot \frac{L}{2} = 1 \Rightarrow A = \sqrt{\frac{2}{L}}$$

$$b) \int_{-\frac{L}{4}}^{\frac{L}{4}} \frac{2}{L} \cdot \left( \frac{1}{2} + \frac{\cos\left(\frac{6\pi x}{L}\right)}{2} \right) dx = \frac{1}{2} + \frac{2}{L} \left( \frac{x}{2} \sin\left(\frac{6\pi x}{L}\right) \right) \Big|_{-\frac{L}{4}}^{\frac{L}{4}}$$

$$= \frac{1}{2} + \frac{1}{6\pi} \left[ \sin\left(\frac{3\pi}{2}\right) - \sin\left(-\frac{3\pi}{2}\right) \right] = \frac{1}{2} - \frac{1}{3\pi}$$

$$c) \frac{\partial \psi}{\partial x} = A \frac{3\pi}{L} \left( -\sin\left(\frac{3\pi x}{L}\right) \right)$$

$$\frac{\partial^2 \psi}{\partial x^2} = -A \left( \frac{3\pi}{L} \right)^2 \cos\left(\frac{3\pi x}{L}\right) = -\left( \frac{3\pi}{L} \right)^2 \psi(x)$$

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$$\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} = E \psi(x) \Rightarrow \frac{\hbar^2}{2m} \left( \frac{3\pi}{L} \right)^2 = E$$

$$d) \quad k^2 = \frac{2mE}{\hbar^2} \Rightarrow k = \frac{n\pi}{L} \Rightarrow E = \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 n^2 \pi^2}{2mL^2}$$

$$E_0 = \frac{\hbar^2 \pi^2}{2mL^2} \Rightarrow \Delta E = E_2 - E_0 = \left( \frac{\hbar^2 \pi^2}{2mL^2} \right) [4 - 1] = \frac{3\hbar^2 \pi^2}{2mL^2}$$

$$E = \frac{hc}{\lambda} \Rightarrow \lambda = hc \cdot \frac{mL}{4\hbar^2 \pi^2} \Rightarrow \lambda = \frac{mLc}{2\pi\hbar}$$

$$Q4. \quad \frac{dN(t)}{dt} = -RN(t)$$

$$a) \quad \frac{dN(t)}{N(t)} = -R dt \Rightarrow N(t) = N(0) \cdot e^{-Rt}$$

$$\frac{N(t)}{N(0)} = \frac{1}{e} = e^{-2 \cdot 10^6 R} \Rightarrow 0 = -2 \cdot 10^6 R + L$$

$$R = \frac{1}{2} \cdot 10^6 \cdot \frac{L}{2}$$

b) No referência de tempo.

$$\bar{t} = \gamma \bar{t} = 35 \cdot 10^6 \text{ s}$$

$$d = \bar{t} \cdot v_m = 35 \cdot 10^6 \cdot 0,999c \approx 1000 \text{ m}$$

$$\therefore N(\bar{t}) = \frac{N_0}{e}$$

$$b) \quad d = \bar{t} \cdot v_m = 2 \cdot 10^6 \cdot 0,999c = 600 \text{ m}$$

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Q5.

a)  $E = \frac{p^2}{2m} + mgh$

$$Z = \int d^3p d^3x e^{-\beta(\frac{p^2}{2m} + mgh)} = \int_0^L dx e^{-\beta mgh} \cdot \int_{-\infty}^{\infty} d^3p e^{-\frac{\beta p^2}{2m}}$$

$$Z = \left(\frac{2m\pi}{\beta}\right)^{\frac{3}{2}} \frac{(e^{-\beta mgh} - 1)}{-\beta mgh}$$

$$\left(\frac{2\pi}{\alpha}\right)^{\frac{3}{2}} \cdot \left(\frac{2m\pi}{\beta}\right)^{\frac{1}{2}}$$

$\alpha = \frac{\beta}{2m}$

11) + dklkl,

$Z = \iiint_0^L \iiint_{-\infty}^{\infty} e^{-\beta(\frac{p_x^2 + p_y^2 + p_z^2}{2m} + mgh)} dp_x dp_y dp_z dz$

$$Z = \underbrace{\left(\frac{2m\pi}{\beta}\right)^{\frac{3}{2}}}_{L^2} \int_0^L e^{-\beta mgh} dz$$

$$Z = \left(\frac{2m\pi}{\beta}\right)^{\frac{3}{2}} \cdot \left[ \frac{1 - e^{-\beta mgh}}{\beta mgh} \right]$$

b)  $\langle u \rangle = -\frac{\partial \ln Z}{\partial \beta} = \frac{3}{2} \frac{2}{\beta} k_B = \frac{3}{2} k_B T$

$$\ln Z = \frac{3}{2} \left[ \ln(2m\pi)^2 - \ln \beta \right]$$

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$$c) \quad \frac{\partial \ln Z}{\partial \beta} = \frac{-\frac{\partial}{\partial \beta} \cdot mgl}{1 - e^{-\beta mgl}} + \frac{1}{\beta}$$

$$\ln Z = \ln(1 - e^{-\beta mgl}) - \ln(\beta mgl)$$

$$= \frac{mgl}{1 - e^{-\beta mgl}} + \frac{1}{\beta}$$

$$d) \quad \frac{\partial}{\partial \beta} \ln Z = \frac{mgl}{1 - e^{-\beta mgl}} + \frac{1}{\beta} = 0$$

$$= \frac{mgl}{1 - e^{-\beta mgl}} + \frac{1}{\beta} = 0$$

06.

$$a) \quad \tau = I \cdot \alpha = \frac{m D^2}{4} \cdot \frac{d\omega}{dt}$$

$$\alpha = \frac{d\omega}{dt}$$

$$\tau = dF = d \cdot \frac{m \omega^2 D r}{2}$$

$$I = \int r^2 dm = \int_{\frac{D}{2}}^{\frac{D}{2}} r^2 \cdot p \cdot dr = p \cdot \frac{r^3}{3} \Big|_{\frac{D}{2}}^{\frac{D}{2}} = \frac{p D^3}{12} = \frac{D^2}{12}$$



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Q8.  $|\psi(0)\rangle = \sum_n c_n |p_n\rangle$

a)  $\sum_n |c_n|^2 = 1$

$|\psi(t)\rangle = \sum_n c_n |\phi_n\rangle \cdot e^{-iE_n t/\hbar}$

Prob  $n=1 =$

$= \sum_{n=2} |c_n|^2 = |c_0|^2 + |c_1|^2$

b)  $|\psi(0)\rangle = c_0 |\phi_0\rangle + c_1 |\phi_1\rangle$

$\langle H \rangle = \langle \psi(0) | \hat{H} | \psi(0) \rangle = (\langle \phi_0 | c_0 + \langle \phi_1 | c_1) \hat{H} (c_0 |\phi_0\rangle + c_1 |\phi_1\rangle)$

$= E_0 c_0^2 + E_1 c_1^2$

$E_0 = \frac{\hbar \omega}{2}, E_1 = \frac{3\hbar \omega}{2}$

$E_0 c_0 |\phi_0\rangle + E_1 c_1 |\phi_1\rangle$

$\langle H \rangle = \hbar \omega \Rightarrow E_0 |c_0|^2 + E_1 |c_1|^2 = \hbar \omega \Rightarrow \frac{|c_0|^2}{2} + \frac{3|c_1|^2}{2} = 1$

$|c_0|^2 + |c_1|^2 = 1$

$|c_0|^2 = 1 - |c_1|^2$

$|c_0|^2 = \frac{3}{4}$

[7]

$\frac{3|c_1|^2}{2} = 1 - \frac{(1 - |c_1|^2)}{2}$

$3|c_1|^2 = 2 - 1 + |c_1|^2 \Rightarrow |c_1|^2 = \frac{1}{4}$

a)  $\vec{r} = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2) = \frac{1}{2} m [\dot{r}^2 + r^2 \dot{\phi}^2] = L$   
 $x = r \cos(\omega t) = \dot{r} \cos(\omega t) - r \sin(\omega t) \cdot \omega$   
 $y = r \sin(\omega t)$

b)  $L = \frac{1}{2} m r^2 \dot{\phi}^2 = 0 \Rightarrow \dot{\phi} = \omega$

$p_\phi = \frac{\partial L}{\partial \dot{\phi}}$

c)  $r = e^{i\omega t}$

$p_r = \frac{\partial L}{\partial \dot{r}} = m \dot{r}$

$\dot{r} = \frac{p_r}{m}$

d)  $H = \sum p_i \dot{q}_i - L$

$H = m \dot{r}^2 - \frac{1}{2} m \dot{r}^2 - r^2 \omega^2 = \frac{1}{2} m \dot{r}^2 - r^2 \omega^2$

$H = \frac{1}{2} \frac{p_r^2}{m} - r^2 \omega^2, \quad \dot{r} = \frac{\partial H}{\partial p_r} = \frac{p_r}{m}, \quad \dot{r} = \frac{\partial H}{\partial r} = 2r\omega^2$

e)  $\frac{\partial L}{\partial t} = 0$  energy conserved

$\frac{\partial L}{\partial \phi} = 0$  Angular momentum conserved

$$c) \langle \psi(0) | \hat{x} | \psi(0) \rangle = (\langle \phi_0 | c_0 + \langle \phi_1 | c_1) \hat{x} (c_0 | \phi_0 \rangle + c_1 | \phi_1 \rangle) = \langle \hat{x} \rangle$$

$$(\langle \phi_0 | c_0 + \langle \phi_1 | c_1) \cdot \left[ \left( c_0 \sqrt{\frac{\hbar}{2m\omega}} \sqrt{0+1} | \phi_1 \rangle + c_1 \sqrt{\frac{\hbar}{2m\omega}} \sqrt{1} | \phi_0 \rangle \right) \right]$$

$$c_0 = |c_0| e^{i\theta_0}$$

$$c_1 = |c_1| e^{i\theta_1}$$

$$c_0 c_1^* \sqrt{\frac{\hbar}{2m\omega}} + c_0 c_1^* \sqrt{\frac{\hbar}{2m\omega}} = \frac{1}{2} \sqrt{\frac{\hbar}{m\omega}}$$

$$c_0 |c_1| e^{i\theta_0} \sqrt{\frac{\hbar}{2m\omega}} + c_1 |c_1| e^{-i\theta_1} \sqrt{\frac{\hbar}{2m\omega}} = \frac{1}{2}$$

$$|c_0| = \frac{\sqrt{3}}{2} \quad |c_1| = \frac{1}{2}$$

$$\theta_0 - \theta_1 = \frac{\pi}{4}$$

$$\sqrt{\frac{\hbar}{2m\omega}} e^{i\theta_0} + e^{-i\theta_1} \sqrt{\frac{\hbar}{2m\omega}} = \frac{1}{2} \Rightarrow e^{i\theta_0} + e^{-i\theta_1} = \frac{\sqrt{2}}{2} e^{i\theta_1}$$

$$\omega(\theta_1) = \sqrt{\frac{2}{3}} \Rightarrow \theta_1 = \omega^{-1}\left(\sqrt{\frac{2}{3}}\right) \quad 30^\circ$$

$$d) |\psi(0)\rangle = \frac{\sqrt{3}}{2} |\phi_0\rangle + \frac{1}{2} e^{i\theta_0} |\phi_1\rangle$$

$$|\psi(t)\rangle = \frac{\sqrt{3}}{2} e^{i\omega_0 t} |\phi_0\rangle + \frac{1}{2} e^{i(\omega_1 t + \theta_0)} |\phi_1\rangle$$

$$\langle \hat{x} \rangle(t) = \langle \psi(t) | \hat{x} | \psi(t) \rangle =$$



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Q9.

$\epsilon_{ijk} = 1$  or  $-1$  or  $0$   
if cyclic index  
if repeated index

a)  $[L_i, L_j] = i\hbar \epsilon_{ijk} L_k$

$$\vec{L} \times \vec{L} = \begin{pmatrix} i & j & k \\ L_i L_j L_k \\ L_i L_j L_k \end{pmatrix} = (L_j L_k - L_k L_j) \hat{i} + (L_k L_i - L_i L_k) \hat{j} + (L_i L_j - L_j L_i) \hat{k}$$

$$[L_j, L_k] \hat{i} + [L_k, L_i] \hat{j} + [L_i, L_j] \hat{k}$$

$$= i\hbar \left[ \underbrace{\sum_i \epsilon_{ijk} L_i}_{\epsilon_{jki} L_i} + \sum_j \epsilon_{ijk} L_j + \sum_k \epsilon_{ijk} L_k \right] = i\hbar [L_i + L_j + L_k] = i\hbar \vec{L}$$

b)  $\vec{L} = R \times P, [R_i, P_j] = i\hbar \delta_{ij}$

$$[L_i, R_j] = L_i R_j - R_j L_i = (R_i \times P_i) R_j - R_j (R_i \times P_i)$$

$$L = R \times P = [R_j, P_k] \hat{i} + [R_k, P_i] \hat{j} + [R_i, P_j] \hat{k}$$

$$L = i\hbar [\delta_{jk} \hat{i} + \delta_{ki} \hat{j} + \delta_{ij} \hat{k}]$$

06.

a)  $P = \frac{aT^4}{3}$ ,  $U = aT^4 V$ ,  $T_0 V_0$   
 $\Delta U = aT_0^4 V_0$

$dU = dQ - dW$   
 $\Delta U = Q - W$

$W = P \Delta V$

$W = \frac{aT_0^4}{3} \cdot V_0$  *van der Waals*

b)  $Q = aT_0^4 V_0 + \frac{aT_1^4}{3} V_0 = \frac{4}{3} T_0^4 V_0$

c)  $0 = \left( \frac{\partial U}{\partial T} \right)_V dT + \left[ \left( \frac{\partial U}{\partial V} \right)_T + P \right] dV$

$0 = 4aT^3 V dT + \left[ aT^4 + \frac{aT^4}{3} \right] dV \Rightarrow 4aT^3 V dT = -\frac{4}{3} aT^4 dV$

d)  $W = -\Delta U = -a \Delta T \Delta V = a(T_1^4 - T_0^4) \cdot V_0$   
 $\frac{-3dT}{T} = \frac{dV}{V} \Rightarrow -3 \ln T = \ln V$

$W = a \left( \frac{T_0^{12}}{16} - T_0^4 \right) V_0 = \frac{aV_0 T_0^{12}}{16} - aT_0^4 V_0$

(10)

$T^3 = V$   
 $V = \frac{1}{T^3} = VT^3 = 1$